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# The Crystal Structure of 5-Halogeno-2-phthalimidobenzoic Acid Monohydrate

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The crystal structures of the isomorphous 5-iodo- and 5-bromo-2-phthalimidobenzoic acid monohydrates ( $C_{15}H_8NO_4I.H_2O \& C_{15}H_8NO_4Br.H_2O$ ) have been determined by two-dimensional Fourier methods. The space group is PI with two molecules in the unit cell. The use of generalized and modulus projections based on  $hk\bar{1}$  reflexions led to the solution of the structure. A suspected ring closure was found not to occur. In addition, the molecule of water, not initially revealed by chemical analysis, was identified.

## Introduction

From certain chemical experiments (Lamchen, 1961), the accepted structure (A) of 2-phthalimidobenzoic acid was considered unlikely. The alternative structures (B) and especially (C) were thought to be more in accordance with the results obtained. To allow a direct structure determination to be undertaken, both an iodine and a bromine derivative were prepared in which the halogen atom had most probably been substituted at the position indicated.



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# Experimental

Crystals of the two compounds were obtained in the form of colourless needles from dilute solutions in ethyl alcohol. X-ray oscillation and Weissenberg photographs showed them to be triclinic with the shortest crystal axis parallel to the needle axis. The unit cells chosen were such that a and b had the shortest translations in real space. The directions of the a, b and c axes were chosen so as to make the cell angles in real space obtuse. This fixed the unit cells (Table 1) uniquely.

Table 1. Crystal data  $C_{15}H_8NO_4Br.H_2O$  $C_{15}H_8NO_4I.H_2O$ 11·06±0·05 Å  $10.91 \pm 0.05$  Å a  $11.29 \pm 0.05$  Å  $11.15 \pm 0.05$  Å b  $6.42\pm0.03$  Å  $6.42 \pm 0.03$  Å с 104° 55′ ± 20′ 106° 24′ + 30′ α 105° 49′ ± 20′  $102^{\circ} 35' \pm 30'$ β 97° 37' ± 20' 97° 30′ ± 30′ γ  $D_m$  $1.86 \pm 0.02 \text{ g.cm}^{-3}$  $1.70 \pm 0.02$  g.cm<sup>-3</sup>  $D_x$ 1.88 g.cm<sup>-3</sup> 1.69 g.cm<sup>-3</sup>

Since no computer was available when this work was started, it was necessary to proceed by twodimensional methods.

Crystals, close to the optimum size, having the following cross-sections were used:

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The hk0 intensities for both derivatives were measured from multiple-film Weissenberg photographs with Cu  $K\alpha$  radiation. Because the shape of the spots on the bromine films was poor, the analysis was concentrated on the iodine derivative. The  $hk\bar{1}$  intensities for this compound were obtained from equi-inclination Weissenberg photographs.

The intensities were estimated visually. For the first-layer spectra, only the peak intensities of the elongated spots were measured and subsequently a correction for the elongation (a geometrical factor) was made (Phillips, 1956). For this purpose, two sets of films were needed to cover the reciprocal-lattice section. The number of observed intensities and the total number of intensities within the limiting sphere for each layer are given in Table 2.

## Table 2. Observed intensities

Layer	$n_{\rm obs}$	$n_{\rm total}$
Br(hk0)	212	293
I(hk0)	198	290
I $(hk\overline{1})$	496	590

Since the linear absorption coefficients were high  $(\mu_{\rm Br}=44 \text{ cm}^{-1} \text{ and } \mu_{\rm I}=193 \text{ cm}^{-1})$ , absorption corrections were applied to the intensities. Albrecht's (1939) method was used and the graphical aids suggested by Rogers & Moffet (1956) speeded up the calculations considerably. The correction factor was calculated as a function of the rotation angle for several values of  $\xi = 2 \sin \theta$ , and the results were contoured on the actual Weissenberg photograph scale. The ratios of the maximum to the minimum correction factor were 1.7 and 1.4 for the iodine and bromine derivatives respectively. Because the equi-inclination angle was small ( $\mu = 7^{\circ}$ ), the absorption corrections did not vary rapidly with  $\mu$  and it was therefore assumed that the correction factors for the first and zero layers were the same.

Wilson plots were used to put the intensities on an approximately absolute scale. A statistical survey of the intensities gave an N(z) distribution favouring the theoretical curve for a centrosymmetric structure (Howells, Phillips & Rogers, 1950), especially when allowance was made for the predominant heavy atoms (Sim, 1958). Since there were two molecules in the unit cell, it was assumed that the space group was  $P\overline{1}$ .

In the structure-factor calculations, the atomic form factors employed were those of Berghuis, Haanappel, Potters, Loopstra, MacGillavry & Veenendaal (1955) for carbon, nitrogen and oxygen and those of Thomas & Umeda (1957) for bromine and iodine.

# Structure determination

### hk0 projections

The hk0 Patterson projections confirmed that the two compounds were isomorphous, the halogen atom having (x, y) coordinates equal to (0.125, 0.116). Isomorphous replacement was used to determine the phases in the first instance. Since the halogen atom was almost in the position  $x \approx y \approx \frac{1}{8}$ , its contribution to some structure factors was small. Only 70% of the signs could be determined with any certainty and the corresponding electron-density projections showed too much false detail to be interpreted.

A survey of the hk0 Patterson function was then undertaken to see if any additional information could be obtained. This included the use of the Buerger minimum function (Buerger, 1959) and the difference Patterson function. The modified Patterson projection (White & Clews, 1956) proved to be more successful. Iodine  $F^2(hk0)$  terms having  $\sin \theta < 0.3$  were omitted to strengthen the scattering power of the iodine atom relative to the light atoms. A Patterson projection, computed with the use of the remaining terms, was considerably sharper than the normal Patterson projection. The modified Patterson projection was vector shifted and the Buerger minimum function drawn. The resulting map was an improvement on the initial projection and gave some indication of possible atom sites.

## Indine $hk\overline{1}$ projection

Generalized and modulus projections used in conjunction have much greater power than the normal zero-layer projections because of the possibility of obtaining all three coordinates (Buerger, 1960). The resolving power of the generalized electron-density projection is also usually greater (Phillips, 1954).

The components of the generalized Patterson projection,  ${}_{c}P_{\overline{1}}$  and  ${}_{s}P_{\overline{1}}$ , defined by

$${}_{c}P_{\bar{1}} = \frac{1}{A} \sum_{hk} F^{2}(hk\bar{1}) \cos 2\pi (hx + ky)$$
$${}_{s}P_{\bar{1}} = -\frac{1}{A} \sum_{hk} F^{2}(hk\bar{1}) \sin 2\pi (hx + ky)$$

gave the z coordinate of the iodine atom as (0.317). The modulus Patterson projection

$$|P_{\bar{1}}| = \frac{1}{(cP_{\bar{1}}^2 + sP_{\bar{1}}^2)}$$

resembled the modified, rather than the normal, hk0 Patterson projection.

The heavy-atom method was used to allocate the signs of  $F(hk\bar{1})$ . Owing to the addition of the third coordinate of the heavy atom, no less than 90% of the phases could be determined, of which 95% proved to be correct. The components of the complex generalized electron-density projection

$$\varrho_{\bar{1}} = c \varrho_{\bar{1}} + i_{s} \varrho_{\bar{1}}$$

were then calculated. From these, the modulus projection  $|\varrho_{\rm I}|$  was obtained, which was more promising than the corresponding hk0 projection. With the aid of the component projections, a model was made of all the likely atomic sites and the approximate position of the molecule was derived.

#### Initial refinement

Initial refinement began on the hk0 data using Fourier and difference Fourier syntheses. Structure factors were calculated by the use of 3° Beevers-Lipson strips (Beevers & Lipson, 1951). Successive syntheses indicated that individual temperature factors were necessary, viz. B=3 Å<sup>2</sup> for the light atoms and B=5 Å<sup>2</sup> for the heavy atom.

The third difference Fourier synthesis showed the position of the oxygen atoms to be such that they could not close the ring in structure (C). This structure could not therefore be the correct one. The atoms were then placed according to the available evidence without reference to any particular assumed structure. The residual  $R = \Sigma ||F_o| - |F_c|| / \Sigma |F_o|$  was 25% for the iodine compound and 36% for the bromine compound. The corresponding Fourier and difference Fourier syntheses both suggested that a molecule of water was also present. This was later confirmed by chemical analysis of the crystals used in the investigation; the initial chemical analysis had been carried out, however, on dried crystals.

## Final refinement

The final stages of the refinement were carried out on the STANTEC ZEBRA computer of the C.S.I.R. using the diagonal approximation of the least squares (Schoone, ZK 22). Successive cycles reduced the residual factor to 15% for the iodine data and to 25% for the bromine data. Difference syntheses suggested considerable anisotropic motion of both halogen atoms. Splitting the heavy atom into two atoms of half weight placed 0.2 Å apart (Vos & Smits, 1961), reduced each R value by another 3%. One refinement cycle, including the temperature factors, reduced the residuals to 9.6% and 17% for the iodine and bromine structure factors respectively.

The resulting (x, y) coordinates were used to calculate approximate z coordinates. First, the two sets of  $(hk\overline{1})$  reflexions were put on a common scale by

Tabl	le 3.	Atomic	parameters.	for	iodine	derivative
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Atom	x/a	y/b	z/c	B
Cl	0.8875	0.2517	0.733	$4 \cdot 0 \text{ Å}^2$
C2	0.7632	0.5009	0.595	$3 \cdot 5$
C3	0.7508	0.1222	0.372	5.3
C4	0.8496	0.1039	0.295	$5 \cdot 0$
C5	0.9753	0.1459	0.467	$4 \cdot 3$
C6	0.9974	0.2260	0.672	4.4
$\mathbf{C7}$	0.9147	0.3273	0.955	$3 \cdot 6$
C8	0.5900	0.3152	0.627	$3 \cdot 5$
C9	0.4906	0.3159	0.760	$4 \cdot 0$
C10	0.4071	0.3927	0.782	<b>4</b> ·0
C11	0.3306	0.3735	0.910	<b>4</b> ·4
C12	0.3370	0.2814	1.002	<b>4</b> ·8
C13	0.4160	0.1986	0.977	$4 \cdot 9$
C14	0.4915	0.2182	0.840	$4 \cdot 6$
C15	0.5939	0.1532	0.750	$3 \cdot 9$
N16	0.6546	0.2184	0.671	4·1
017	0.6275	0.0636	0.830	<b>4</b> ·7
018	0.6260	0.3950	0.561	4.4
O19	0.8200	0.3593	1.024	$4 \cdot 6$
O20	1.0263	0.3663	1.064	4.9
O21	0.8525	0.4898	0.472	4.6
122	1.1308	0.0999	0.3364	5.4
$\mathbf{I22}$	1.1397	0.1219	0.3380	5.4

calculating structure factors for each set and using the scale factors so obtained. Successive refinements of the positional parameters reduced R to 16%. From a comparison between the Fourier and leastsquares z coordinates, a final set of z coordinates was adopted. With one refinement cycle of the temperature factors, this set of coordinates (Table 3) gave a final residual of 13% (observed reflexions only).

### Discussion

The atomic maxima of the iodine  $hk\overline{1}$  modulus projection (Fig. 1) were more clearly resolved than in the corresponding hk0 projection owing in part to the larger number of reflexions with high  $\sin \theta$ . Final difference syntheses on the hk0 data were clean, showing that all the atoms had been correctly placed. Splitting the heavy atom into two did not compensate completely for its anisotropic vibration.



Fig. 1. Iodine  $hk\overline{1}$  modulus electron-density projection. Contour intervals=1 e.Å<sup>-2</sup> starting at 2 e.Å<sup>-2</sup>, except that for I atom contour intervals=2 e.Å<sup>-2</sup> and 10 e.Å<sup>-2</sup> above 10 e.Å<sup>-2</sup>.



Fig. 2. Inter-atomic distances (Å) of iodine derivative.

The interatomic distances (Fig. 2) are within the range expected, taking into account the inaccuracies in the determination of the z coordinates. Cruick-shank's (1960) formula gives the estimated standard

Table 4. Direction cosines of best planes with respect to axes  $\mathbf{a}^*$ ,  $\mathbf{c} \wedge \mathbf{a}^*$  and  $\mathbf{c}$ 

#### Plane 1 contains atoms 1-7, 16, 19, 20 and 22 Plane 2 contains atoms 8-18

	Plane 1	Plane 2
L	+0.010	+0.693
M	+0.813	+0.445
N	-0.582	+0.568
D	-0.131	+ 5.989

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Table 5. Observed and calculated structure factors for iodine derivative To put  $F_o$  and  $F_c$  on an absolute scale, divide by 5

ĸ	Fo	Fc	ĸ	Fo	Fc	ĸ	Fo	Fc	K	Fo	Fc	ĸ	Fo	Fc	ĸ	Fo 11K1	Fc	ĸ	Fo	Fc	к	Fo	Fc
12345 <b>6</b> 789	050 271 214 190 465 197 91 53 126 107	278 -234 -180 -453 -222 - 90 - 80 138 93	0 1 2 3 4 5 6 7 8	5400 18 24 239 149 71 33 80 75 42	00 30 153 53 35 - 97 - 75 - 32	- 3 - 4 - 7 - 8 - 9	142 78 71 71 27 11 <b>50</b> 91	152 85 - 61 - 63 - 19 - 76	9 10 11 - 1 - 2 - 3 - 4 - 5	69 68 35 15 41 67 256 198 56	85 63 44 - 32 - 8 - 56 -280 -256 - 60	5 6 8 9 10 - 1 - 2 - 3 - 3 - 4	153 45 35 40 34 125 20 217 139	133 35 - 38 - 42 - 35 -115 - 40 193 102	0 2 3 4 5 6 - 1 - 2 - 3	69 34 75 58 17 8 62 107 56	51 - 26 - 46 - 42 - 31 12 53 79 52	5 6 7 8 10 11 12 - 1 - 2 3	12 184 182 128 27 33 31 70 431	18 200 169 124 - 24 - 28 - 48 110 472 257	12 - 1 - 2 - 3 - 4 - 5 - 6 - 7	24 102 142 81 26 50 56 32	26 -118 -154 - 85 - 41 55 43 32
10 0 1 2 3 4	50 1 <b>60</b> 251 183 85 334 229	45 233 - 92 -341 -281	- 1 - 2 - 3 - 4 - 5 - 6 - 7 - 9 -10	228 296 111 68 293 261 96 57 66	-303 -303 -112 - 85 279 270 84 - 68 - 77	1 - 1 - 3 - 4 - 5 - 8 - 9	27 47 47 75 57 47 47 12 <b>60</b>	- 43 - 38 53 61 49 - 54 - 36	- 6 - 7 - 8 - 9 -10 -12 -13	72 56 142 93 49 24 23 3 <b>%</b> 1	-100 73 121 103 60 - 40 - 26	- 5 - 6 - 7 - 8 - 9 -10 -13	208 89 78 122 96 18 3 7K1	173 83 - 85 - 103 - 89 - 15 12	- 4 - 5 - 6 - 7 - 9 -10	16 87 68 55 24 19 12K1 16	- 24 - 79 - 43 - 49 27 19	- 3 - 4 - 5 - 6 - 7 - 8 -10 -11 -12	230 113 62 116 120 81 44 51 4	257 137 - 42 - 96 -119 - 63 47 37 9	0 1 2 3 4 5 6 7	8K1 128 44 50 133 156 32 64	-142 - 40 53 139 165 21 - 72
5 6 7 8 9 12 - 1 - 2 - 3 - 4	98 136 132 75 27 244 351 18 116	-104 112 125 140 81 - 27 203 339 15 -133	-11 0 1 2 3 5 6	42 6879 139 204 203 121 87 62	- 49 157 197 213 140 - 85 - 80	3 4 - 1 - 2 - 5 - 6	27 19 63 65 19 50 13 <b>KO</b>	22 18 - 56 - 58 37 47	0 1 2 3 4 5 6 7 8	477 69 28 258 149 268 49 102 95	519 71 - 23 -321 -194 -287 - 60 120 109	0 2 3 4 5 6 7 8 9	264 88 116 88 106 21 49 57 35	-251 103 145 93 116 - 16 - 43 - 57 - 41	1 2 3 5 1 2 3 5 1 2 3 5 1 2 3 4 6 7	75 49 16 7 24 49 68 20 48 58	- 63 - 41 - 3 22 22 34 53 15 - 38	0 1 2 3 4 5 6 7	481 132 44 155 178 242 165 123 111	163 23 -175 -163 -221 -159 127 77	8 9 12 - 1 - 2 - 3 - 4 - 5 - 6	51 75 44 23 71 30 37 44 84 39	- 40 - 81 - 40 28 - 62 - 44 - 34 55 76 32
- 6 - 7 - 8 - 9 -10 -11 -12	202 232 135 85 76 38 60 27 2 <b>30</b>	-313 -209 -142 112 83 <b>21</b> 62 34	- 1 - 2 - 3 - 4 - 6 - 7 - 8 - 9 -11	65 163 231 214 107 52 73 125 50 27	- 74 -168 -237 -244 -114 54 65 130 53 - 32	- 1 - 2 - 3 0 1 2 3	19 27 27 0%1 179 388 325 938	- 18 - 29 - 28 -219 -357 -309 818	11 13 - 1 - 2 - 3 - 4 - 5 - 6 - 7	13 11 580 295 105 232 354 117 69	39 - 14 - 15 507 251 87 -230 -360 - 96 - 86	- 1 - 2 - 3 - 4 - 5 - 6 - 7 - 9 -10	288 195 64 59 263 90 68 30 47	-246 -155 - 62 58 228 74 56 - 25 - 48	- 8 - 9 0 1 2 - 2	30 8 13KÎ 55 35 20 18	- 12 - 8 - 44 - 24 - 26 - 15	8 9 11 12 13 - 1 - 2 - 3 - 4 - 5	234 120 48 34 23 96 299 9 73 114	208 118 - 56 - 31 - 35 117 303 - 18 - 92 -106	0 1 2 3 4 5 6	9K1 82 62 51 39 47 20 90	- 82 - 75 - 65 55 70 62
0 1 2 3 4 5 6 7 8	50 182 599 225 60 95 149 135 22	67 -251 -599 -282 - 62 106 169 126 23	0 1 2 3 4 5 6	7120 175 69 116 68 126 95 42	169 92 114 - 64 -108 - 99 - 38	4 5 7 8 9 10 11 12 13	204 29 210 217 141 42 11 69 62 13	235 27 -207 -247 -151 - 37 - 35 81 66 14	- 8 - 9 -10 -11 -13 0 1	49 74 37 34 11 4 <b>81</b> 178 262	60 59 34 47 - 20 159 -271	-11 -12 0 1 2 3 4	35 13 8K1 49 144 110 29	- 31 - 24 - 89 62 144 119 37	- 3 - 4 - 5 - 6 - 7 0 1	41 45 9 8 5 1K1 512 96	35 36 8 15 - 3 -568 -104	- 6 - 7 -10 -11	114 20 53 2 5K1 198 267	-110 - 2 51 22 241 262	7 9 10 11 12 - 1 - 2 - 3 - 4 - 5	73 43 32 16 15 72 21 37 34 18	59 - 45 - 34 - 18 1 - 70 31 25 35 21
$ \begin{array}{r} 10 \\ 11 \\ -2 \\ -3 \\ -4 \\ -6 \\ -7 \\ -8 \\ -10 \\ \end{array} $	53 42 114 490 477 1 <b>60</b> 109 158 97 42	- 60 - 48 104 393 470 149 -128 -163 - 95	8 - 1 - 3 - 4 - 5 - 7 - 7 - 7 - 7 - 9	50 154 197 161 164 71 107 66	32 163 -204 -147 -149 67 131 61	- 1 - 2 - 3 - 4 - 5 - 6 - 7 - 8 - 9	164 512 65 45 253 262 113 09 42	-152 -516 - 49 - 36 266 274 87 5 - 61	234 5 7 8 9 - 1	135 144 87 38 68 100 108 22 134	-179 -204 -122 55 67 94 125 7 133	5 6 7 8 - 1 - 2 - 3 - 4 - 5	42 52 25 39 204 125 23 21	- 50 - 45 - 43 - 31 - 28 -155 -117 - 26 24	2 3 4 5 6 7 8 9 10	116 673 64 192 151 32 92 126 57 30	101 585 67 197 133 - 62 - 77 -127 - 53 62	2 3 4 5 6 8 9 10 13 - 1	110 151 183 242 143 60 105 56 22 137	113 -147 -178 -219 -124 56 102 58 - 40 172	- 6 1 2 3 5 6 7	15 10K1 41 56 63 39 76 50	20 - 55 - 54 - 63 - 40 - 87
0 1 2 4 5 6	42 3 <b>80</b> 204 294 292 148 242 96	-241 -311 -301 167 296 60	0 1 2 3 4 5 8 9	61 D 190 65 38 112 57 42 33 19	197 68 - 31 - 96 - 50 - 52 37 24	-10 -11 -12 0 1 2 3	50 22 24 1K1 173 277 243 603	- 62 - 10 - 44 160 327 294 624	- 2 - 3 - 4 - 5 - 6 - 7 - 8 - 9 -10 -11	390 235 54 156 247 218 42 34 72 51	-343 206 41 -146 -210 -204 - 33 41 73 49	- 6 - 7 - 8 -10 -11 -12	159 153 23 27 43 24 9K1 128	140 141 8 - 39 - 44 - 33	13 - 1 - 2 - 3 - 4 - 5 - 6 - 8 - 9	14 763 44 167 118 232 110 74 67	17 -852 - 54 200 131 237 112 - 49 - 35	- 23 - 34 - 56 - 89 - 9	17 74 115 63 59 29 641	21 - 75 -136 - 65 - 62 59 21	9 10 11 - 1 - 2 - 3 4	11 24 13 15 67 32 24 11K1	49 - 16 - 32 - 28 11 64 34 14
7 9 10 11 - 1 - 2 - 3 - 4 - 5 - 7	71 75 42 131 114 424 124 229 123	77 - 70 - 48 - 46 -121 113 370 153 248 -136	- 1 - 2 - 3 - 4 - 5 - 6 - 7 - 9 -10	170 67 33 98 111 91 47 53 63	173 71 - 51 - 98 -106 - 72 - 58 48 64	4 5 6 7 8 9 10 11 12	120 196 188 124 84 25 59 34 25 9	110 -191 -185 -146 - 93 33 47 23 20 7	-12 0 1 2 3 4 5	28 5K1 223 286 224 146 24 170	43 -236 -330 -281 -164 30 165	1 2 3 4 5 6 7 8 9	56 90 39 21 49 74 19 15 10	39 109 44 - 16 - 33 - 68 - 7 23 - 16	-10 -11 -12 0 1 2 3	29 34 19 2K1 119 608 164 121	- 34 - 34 23 -155 -612 -111	0 1 2 3 4 5 6 7 8	75 193 115 168 99 201 185 94 39	86 208 119 198 -102 -202 -173 - 80 67	2 3 4 5 6 7 8 9	47 34 33 32 29 22 15 22	- 64 - 45 - 36 - 45 51 40 27 28 25
- 8 - 9 -10 -11 -12	146 121 38 38 47 480 499	-136 -132 - 43 52 43 -485	2 3 4 8 - 1 - 2 - 3	9 <b>68</b> 82 42 19 85 <b>50</b> 68	- 68 - 63 - 38 - 22 - 94 - 31 - 51	- 1 - 2 - 3 - 4 - 5 - 6 - 7 - 8 - 9	296 427 310 306 121 135 105 55 30	-295 -415 -303 -320 130 118 102 75 40	6 7 9 10 11 - 1 - 2 - 3 - 4	152 48 25 49 33 29 381 402 158	122 39 - 28 - 38 5 45 383 363 135	- 1 - 2 - 3 - 4 - 5 - 7 - 8 - 9 - 10 - 11	10 72 116 154 40 30 66 38 10 5	- 58 -100 -131 - 40 31 52 37 16 - 4	4 5 6 7 8 9 10 11 - 1	132 329 250 83 62 59 106 85 261	129 322 220 52 - 65 - 41 - 93 - 74 298	9 10 11 12 - 1 - 2 - 3 - 4 - 6	32 47 69 9 64 93 114 100 15	43 36 74 5 - 67 - 112 - 120 - 92 23	- 2 - 3 0 3 4 5	39 15 12K1 20 17 34 16	41 15 28 - 14 - 44 - 19
134568012	156 101 160 191 71 90 47 227 356	-164 91 164 206 70 - 74 - 36 -223 -324	- 4 - 5 - 6 - 7 - 8 - 9 -10 -11	100 89 124 57 27 27 33 33	- 49 - 118 - 49 - 27 - 27 - 24 - 33 - 24	-10 -11 -12 -13	24 29 14 8 2K1 361 326	- 39 - 27 - 9 - 5 375 379	- 5 - 6 - 7 - 8 - 9 - 10 -11 - 12	18 41 141 101 63 7 45 21	2 - 17 -149 - 82 - 62 - 62 - 3 - 63 - 25	-12 0 1 3 4 5	7 10 <b>K</b> 1 146 21 48 59 32	- 17 122 - 19 - 60 - 53 - 16	- 2 - 3 - 4 - 5 - 6 - 7 - 7 - 8 - 9	58 146 274 62 17 122 135 67	58 154 246 49 1 -129 -115 - 49	- 7 - 8 1 2 3 4	62 37 7K1 114 136 145 103	61 28 121 142 173 104	8 - 1 1 2 4 5	13 23 13K1 10 11 8 10	33 34 39 28 - 20 - 29
- 3 - 4 - 5 - 6 - 8 - 0 - 10	87 320 255 185 85 38 89	113 297 225 181 - 85 - 19 -109	0 1 2 5 6	33 112 57 38 30	- 37 - 86 - 46 - 40 16	2 3 4 5 6 7 8	31 23 197 200 86 72 69	- 34 - 31 -203 -248 - 95 - 78 85	0 1 2 3 4	6K1 158 238 6! 144 97	-145 -240 - 72 163 113	- 1 - 3 - 4 - 5 - 6 - 8 - 9 -10	197 67 78 138 44 34 88 34	152 - 64 - 46 - 109 - 28 - 31 - 68 - 38	0 1 2 3 4	3K1 86 313 409 426 76	- 99 -323 -380 -350 - 66	5 6 7 8 9 10 11	26 140 96 103 30 20 31	41 -109 -106 -105 - 22 21 48			

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deviation of the light and heavy atom coordinates as 0.015 and 0.006 Å respectively. Since the z coordinates are determined by only one non-zero layer line, the assumption of isotropic sampling is not really valid and the above values can be expected to apply only to the x and y coordinates. From the way in which the z coordinates were determined, the uncertainty in these coordinates seems to be about two to three times that of the other two coordinates.

Using the method of Schomaker, Waser, Marsh & Bergman (1959), the best planes were calculated through the two parts of the molecule. The direction cosines of these planes are given in Table 4. Atoms deviating furthest from their respective planes are (Fig. 2) C4, C5, C15 and O18 (0·10 Å). The two planes have twisted through an angle of 88° with respect to each other about the C–N bond. Since the acid grouping lies in the plane of the benzene ring, a relatively short distance results between N16 and O19 (2·48 Å). This approach is less than the sum of the effective van der Waals radii (2·9 Å), but more than the sum of the effective ionic radii (2·0 Å). This is mainly due to electrostatic attraction between the oppositely charged ions.



Fig. 3. Arrangement of molecules viewed down the b axis. The broken lines represent hydrogen bonds.

Hydrogen bonding (Fig. 3) between the molecule of water and O18 in one molecule (2.85 Å) and O19 in the molecule a - c translation away (2.77 Å) helps bind the molecules within a stack while the stacks themselves are held together by forces of the van der Waals type. The bond angle between the oxygen atoms involved in the hydrogen bonding is 130°.

The individual isotropic temperature factors were obtained by least-squares refinement. Their standard deviation is about 0.5 Å<sup>2</sup>. Atoms on the periphery of the molecule appear to execute thermal motion with a larger amplitude.

In the above analysis, extensive use of generalized

functions has been made. The generalized Patterson projection was sharper than the corresponding zerolayer projection. The determination of the third coordinate, albeit roughly, from the initial generalized projection, allowed a successful trial structure to be found. Lastly, the final generalized projection gave a good check on the least-squares coordinates. The successful application of these methods shows that the scope and power of the generalized projection can, in general, considerably extend the range of molecules that can be solved by two-dimensional methods.

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